# Indian Institute of Science <br> ME 303: Endsemester Exam 

Date: 19/4/2023.
Duration: 9.30 a.m. -12.30 a.m.
Maximum Marks: 100

1. This problem can be treated as a two-dimensional steady-state problem. Consider a circular inclusion of radius $R$, and with conductivity $k_{1}$ embedded in an unbounded domain with conductivity $k_{2}$. The far-field temperature is given by $V_{0}+V_{1} x y$, i.e., $\lim _{r \rightarrow \infty} T=V_{0}+V_{1} x y$, where $r$ denotes the radial coordinate in a polar coordinate system, and $y=r \sin \theta$. Find the steady-state temperature distribution within the inclusion and in the surrounding region. Your solution procedure should be systematic with no guesswork allowed.
2. The governing equation for the velocity component $u_{z}(x, y)$ for Poiseuille flow through a conduit of rectangular cross section (see Fig. 1) is given by

$$
\frac{\partial^{2} u_{z}}{\partial x^{2}}+\frac{\partial^{2} u_{z}}{\partial y^{2}}=-\frac{G}{\mu}
$$

where $G$ and $\mu$ are constants. The boundary condition is the no-slip condition $\left.u_{z}\right|_{x= \pm a / 2}=$ $\left.u_{z}\right|_{y= \pm b / 2}=0$. Find the (steady-state) velocity distribution $u_{z}(x, y)$. Once again, the solution should be derived systematically with no guesswork allowed. State all the assumptions that you make regarding the nature of the solution.
3. An infinitely long circular cylinder of radius $a$ (i.e., the domain is $[0, a] \times(-\infty, \infty)$ in the $r-z$ plane) with insulated boundary $r=a$, is subjected to a heat input $Q(r, z, t)$. Derive the temperature distribution $T(r, z, t)$. You may also directly use the following result for $u=J_{k}(\lambda r), Y_{k}(\lambda r)$ :

$$
\frac{1}{r} \frac{d}{d r}\left(r \frac{d u}{d r}\right)=-\left(\lambda^{2}-\frac{k^{2}}{r^{2}}\right) u
$$



Figure 1: Flow through a rectangular duct.
4. The equation for torsional oscillations of a circular cylinder of radius $a$ and length $L$ is given by

$$
\begin{equation*}
\frac{\partial}{\partial r}\left(\frac{1}{r} \frac{\partial\left(r u_{\theta}\right)}{\partial r}\right)+\frac{\partial^{2} u_{\theta}}{\partial z^{2}}=\frac{1}{c^{2}} \frac{\partial^{2} u_{\theta}}{\partial t^{2}}, \tag{35}
\end{equation*}
$$

where $c$ is the wave speed. With $\rho$ denoting the density, the boundary conditions are $\left.u_{\theta}\right|_{z=0}=0$, and

$$
\left.\frac{\partial u_{\theta}}{\partial z}\right|_{z=L}=\frac{r T(t)}{\rho c^{2}} .
$$

Assume that the boundary conditions on the surface $r=a$ are automatically satisfied.
(a) By substituting $u_{\theta}(r, z, t)=\operatorname{rw}(z, t)$, transform Eqn. (1) into a partial differential equation in terms of $w(z, t)$. Find the boundary conditions also in terms of $w(z, t)$.
(b) Assuming homogeneous initial conditions, find the solution for $w(z, t)$ in terms of $T(t)$.
(c) For the special case $T(t)=T_{0}$, where $T_{0}$ is a constant, find the solution (you need not evaluate summations that arise).
(d) By multiplying the governing equation for $w(z, t)$ that you have derived by an appropriate quantity, and integrating over the domain $[0, L]$, derive an equation of the form

$$
\begin{equation*}
\frac{d H}{d t}=\text { RHS }, \tag{2}
\end{equation*}
$$

where

$$
H=\int_{0}^{L}[\ldots] d z,
$$

where the integrand in the above expression is in terms of $w$ and its derivatives, and the RHS in Eqn. (2) comprises boundary terms involving $T(t)$. If $T(t)$ is suddenly set to zero at say time $t_{0}$, deduce a 'conservation law' for $H$ from Eqn. (2) for $t \geq t_{0}$.

