

Indian Institute of Science

ME 242: Final Exam

Date: 5/12/11.

Duration: 9.30 a.m.–12.30 p.m.

Maximum Marks: 100

1. A circular cylinder of radius a and length L is subjected to a prescribed displacement $u_z = \bar{u}$ at its ends as shown in Fig. 1. Assume that the circular cylinder remains circular and prismatic after deformation. Make an intelligent guess of the form of (u_r, u_θ, u_z) as a function of (r, θ, z) , and use that to find the displacement and stress fields. (Hint: Are the stresses and strains constant?) (20)

2. A truncated wedge of unit thickness (in the z -direction) is loaded by a traction distribution on the top and bottom edges, whose resultant is an axial total load P , as shown in Fig. 2. The lateral edges $\theta = \theta_0$ are traction-free. The stress distribution is given by (30)

$$\begin{aligned}\tau_{rr} &= \frac{c_1 \cos \theta}{r}, \\ \tau_{\theta\theta} &= c_2 \log r + \frac{c_3 \sin \theta}{r}, \\ \tau_{r\theta} &= c_4 r,\end{aligned}$$

with the remaining stresses zero (plane-stress conditions). Treating this problem as two-dimensional, and using *only* the boundary conditions, find the constants c_1 , c_2 , c_3 and c_4 .

3. A truncated hollow sphere with inner and outer radii R_1 and R_2 , is subjected to torsional loading as shown in Fig. 3. The torque, which is generated by a traction distribution on the flat ends, is given by $T\mathbf{e}_z$, and the stress (25)

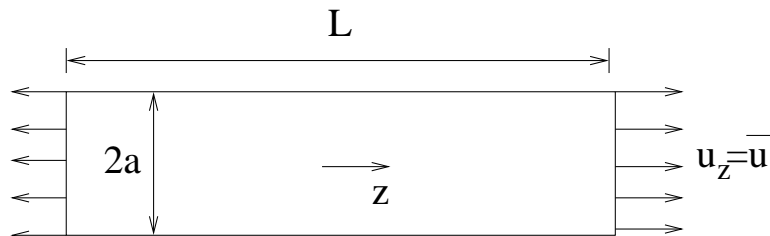


Figure 1: Circular cylinder subjected to a prescribed end displacement $u_z = \bar{u}$

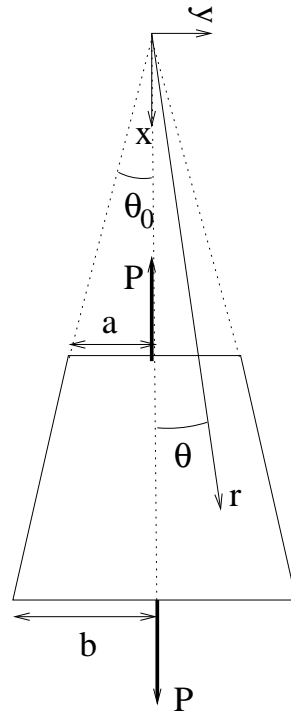


Figure 2: Truncated wedge loaded by an axial load P

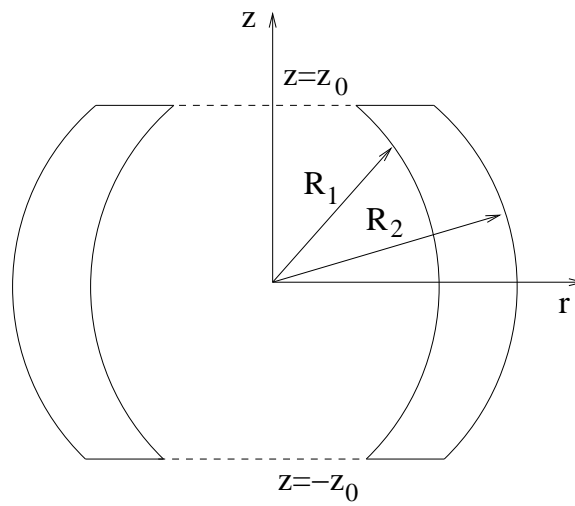


Figure 3: Truncated hollow sphere subjected to torsion.

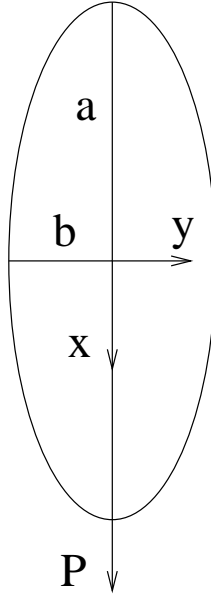


Figure 4: Beam loaded by a terminal load.

distribution is given by

$$\begin{aligned}\tau_{r\theta} &= -\frac{3c}{r^2}z\sqrt{r^2+z^2}, \\ \tau_{\theta z} &= \frac{3c}{r}\sqrt{r^2+z^2},\end{aligned}$$

with the other stresses zero. Find c in terms of T , z_0 , R_1 and R_2 .

4. A beam of length L whose cross-section is given by the equation (25)

$$\left(1 - \frac{x^2}{a^2}\right)^{2\nu} = \left(\frac{y}{b}\right)^2,$$

where ν is the Poisson ratio, is loaded by a traction distribution on the end-face $z = L$ with resultant $P\mathbf{e}_x$, as shown in Fig. 4. The lateral surfaces are traction-free. Assuming

$$\begin{aligned}\tau_{xz} &= c_1(a^2 - x^2), \\ \tau_{yz} &= c_2xy, \\ \tau_{zz} &= -E(L - z)(\kappa_x x + \kappa_y y),\end{aligned}$$

where

$$\kappa_x = \frac{I_{xx}W_x + I_{xy}W_y}{E(I_{xx}I_{yy} - I_{xy}^2)}, \quad \kappa_y = \frac{I_{xy}W_x + I_{yy}W_y}{E(I_{xx}I_{yy} - I_{xy}^2)},$$

and after making appropriate assumptions on the nature of the other stress components, determine c_1 and c_2 using the equations of equilibrium, and the boundary conditions (do not find I_{xx} or I_{yy} explicitly in terms of a and b). State (without proof, but with justification), the values of $\int_A \tau_{xz} dA$, $\int_A \tau_{yz} dA$ and $\int_A (x\tau_{yz} - y\tau_{xz}) dA$.

Some relevant formulae

$$\begin{aligned}
 \epsilon_{rr} &= \frac{\partial u_r}{\partial r}, & \epsilon_{r\theta} &= \frac{1}{2} \left[\frac{1}{r} \frac{\partial u_r}{\partial \theta} + r \frac{\partial}{\partial r} \left(\frac{u_\theta}{r} \right) \right], \\
 \epsilon_{\theta\theta} &= \frac{1}{r} \left(\frac{\partial u_\theta}{\partial \theta} + u_r \right), & \epsilon_{\theta z} &= \frac{1}{2} \left(\frac{\partial u_\theta}{\partial z} + \frac{1}{r} \frac{\partial u_z}{\partial \theta} \right), \\
 \epsilon_{zz} &= \frac{\partial u_z}{\partial z}, & \epsilon_{rz} &= \frac{1}{2} \left(\frac{\partial u_r}{\partial z} + \frac{\partial u_z}{\partial r} \right).
 \end{aligned}$$

$$\boldsymbol{\tau} = \lambda(\text{tr } \boldsymbol{\epsilon})\mathbf{I} + 2\mu\boldsymbol{\epsilon}.$$

$$\begin{aligned}
 \frac{\partial \tau_{rr}}{\partial r} + \frac{1}{r} \frac{\partial \tau_{r\theta}}{\partial \theta} + \frac{\partial \tau_{rz}}{\partial z} + \frac{\tau_{rr} - \tau_{\theta\theta}}{r} &= 0, \\
 \frac{\partial \tau_{\theta r}}{\partial r} + \frac{1}{r} \frac{\partial \tau_{\theta\theta}}{\partial \theta} + \frac{\partial \tau_{\theta z}}{\partial z} + \frac{\tau_{r\theta} + \tau_{\theta r}}{r} &= 0, \\
 \frac{\partial \tau_{zr}}{\partial r} + \frac{1}{r} \frac{\partial \tau_{z\theta}}{\partial \theta} + \frac{\partial \tau_{zz}}{\partial z} + \frac{\tau_{zr}}{r} &= 0.
 \end{aligned}$$

$$[\bar{\mathbf{v}}] = \mathbf{Q}[\mathbf{v}]; \quad [\bar{\boldsymbol{\tau}}] = \mathbf{Q}[\boldsymbol{\tau}]\mathbf{Q}^T, \quad \text{where } Q_{ij} = \bar{\mathbf{e}}_i \cdot \mathbf{e}_j$$

$$\begin{aligned}
 \int_{-a}^a \frac{dy}{(a^2 \cot^2 \theta_0 + y^2)} &= \frac{2\theta_0 \tan \theta_0}{a}, \\
 \int_{-a}^a \frac{dy}{(a^2 \cot^2 \theta_0 + y^2)^2} &= \frac{\tan^3 \theta_0}{a^3} [\theta_0 + \sin \theta_0 \cos \theta_0].
 \end{aligned}$$