

Natural frequency

Hamilton's principle

$$\text{Extremize } u(x, t) \quad A = \int_0^t L dt$$

$$L = \text{Lagrangian} = KE - PE$$

Bar

$$KE = \int_0^L \frac{1}{2} (\rho A dx) \dot{u}^2 = \int_0^L \frac{1}{2} \rho A \dot{u}^2 dx$$

$$\dot{u} = \frac{du}{dt}$$

$$PE = SE + WP$$

= 0 because of free vibration

$$= \int_0^L \frac{1}{2} EA u'^2 dx$$

$$\text{Extremize } u(x, t) \quad A = \int_0^t \int_0^L \underbrace{\left[\frac{1}{2} \rho A \dot{u}^2 - \frac{1}{2} EA u'^2 \right]}_L dx dt$$

$$\frac{\partial L}{\partial u} - \frac{\partial}{\partial x} \left(\frac{\partial L}{\partial u'} \right) - \frac{\partial}{\partial t} \left(\frac{\partial L}{\partial \dot{u}} \right) = 0$$

$$0 - \left(-EAu' \right)' - \frac{\partial}{\partial t} (\rho A \dot{u}) = 0$$

$$(EAu')' - \rho A \ddot{u} = 0$$

$$u(x, t) = \Phi(x) q(t)$$

$$u' = \phi' \dot{q}$$

$$u'' = \phi'' \dot{q}$$

$$\dot{u} = \phi \ddot{q}$$

$$\ddot{u} = \phi \ddot{q}$$

$$\text{Let } A(x) = A_0$$

~~$$EA \phi'' \dot{q}$$~~

$$EA u'' = \rho A \ddot{u}$$

$$E \phi'' \dot{q} = \rho \phi \ddot{q}$$

$$\Rightarrow \frac{E \phi''}{\rho \phi} = \frac{\ddot{q}}{\dot{q}} = \text{Constant} = -\lambda$$

$$E \cdot \ddot{q} + \lambda \rho q = 0$$

$$\boxed{E \phi'' + \lambda \rho \phi = 0}$$

Eigenvalue problem
when A is constant

$$\ddot{q} + \lambda q = 0$$

$$q(t) = B \sin \sqrt{\lambda} t$$

$$\dot{q} = +B \sqrt{\lambda} \cos \sqrt{\lambda} t$$

$$\ddot{q} = -B \lambda \sin \sqrt{\lambda} t$$

$$\text{@ } t=0 \quad \dot{q}=0$$

$$B \sin \sqrt{\lambda} t = 0$$

Since $B \neq 0$

$$\sin \sqrt{\lambda} t = 0$$

$$\text{Natural frequency: } \sqrt{\lambda_n} = n\pi$$

$$n = 1, 2, \dots, \infty$$

$$E \phi_n'' + \lambda_n \rho \phi_n = 0$$

Eigen equation

No space ω
eigen —

$$\sqrt{\lambda} = \omega = \text{Natural freq}$$

ϕ_i = eigenfunction or mode shape

Properties of Eigenfunction

1) ϕ_i 's are orthogonal $\Rightarrow \int_0^L p \phi_i \phi_j dx = 0$

$$\phi_j (E \phi_i'' + p h_i \phi_i = 0)$$

$$\phi_i (E \phi_j'' + p h_j \phi_j = 0)$$

$$E \phi_j \phi_i'' + p h_i \phi_i \phi_j = 0$$

$$E \phi_i \phi_j'' + p h_j \phi_i \phi_j = 0$$

~~$$\phi_j \phi_i'' = \phi_i \phi_j''$$~~

~~$$\frac{\phi_i''}{\phi_i} = \frac{\phi_j''}{\phi_j}$$~~

$$E \phi_j \phi_i'' - E \phi_i \phi_j'' + p (h_i - h_j) \phi_i \phi_j = 0$$

$$\int_0^L (E \phi_j \phi_i'' - E \phi_j'' \phi_i) dx + \int_0^L p (h_i - h_j) \phi_i \phi_j dx = 0$$

~~$$E \phi_j \phi_i' \Big|_0^L - \int_0^L E \phi_j' \phi_i' dx - E \phi_i \phi_j' \Big|_0^L + \int_0^L E \phi_j' \phi_i' dx +$$~~

$$\int_0^L p (h_i - h_j) \phi_i \phi_j dx = 0$$

Boundary Condition
lead to 0

$$\therefore \int_0^L p (h_i - h_j) \phi_i \phi_j dx = 0 \quad \text{when } \underline{i \neq j}$$

When $i=j$

$$2) \int_0^L \rho \phi_i^2 dx = 1 \quad \left[\text{Normalization Condition} \right]$$

$$3) \boxed{\phi(x) = \sum_{i=1}^{\infty} c_i \phi_i(x)}$$

Superposition of mode shapes

$$\cancel{\phi_i} + \dots + \cancel{c_n \rho \phi_i} = 0$$

$$\int_0^L \rho \phi_i dx = \int_0^L \left(\sum_{j=1}^{\infty} c_j \phi_j \right) \rho \phi_i dx$$

$$= \sum_{j=1}^{\infty} \int_0^L \rho c_j \phi_j \phi_i dx = \sum_{j=1}^{\infty} c_j \int_0^L \rho \phi_j^2 dx = c_i$$

$$\boxed{\int_0^L \rho \phi \phi_i dx = \underline{\underline{c_i}}}$$

$$\# \int_0^L \rho \phi^2 dx = 1$$

$$\int_0^L \rho \phi \left[\sum_{i=1}^{\infty} c_i \phi_i(x) \right] dx = 1$$

$$\sum_{i=1}^{\infty} \left(c_i \int_0^L \rho \phi \phi_i dx \right) = 1$$

$$\boxed{\sum_{i=1}^{\infty} c_i^2 = 1}$$

In eigenvalue problem, always ~~$\int_0^L \rho \phi_i dx = 1$~~

$$\int_0^L \rho \phi_i dx = 1$$

$$① E \phi_i'' + \lambda_i \rho \phi_i = 0$$

Equation obtained after substituting $u = \phi(x) R(t)$

$$② E \phi \phi_i'' + \lambda_i \rho \phi \phi_i = 0$$

$$\int_0^L (E \phi \phi_i'' + \lambda_i \rho \phi \phi_i) dx = 0$$

Multiply by ϕ & integrate

$$E \phi \phi_i' \Big|_0^L - \int_0^L E \phi' \phi_i' dx + \lambda_i \int_0^L \rho \phi \phi_i dx = 0$$

Boundary Condition.

$$\lambda_i = \frac{\int_0^L E \phi' \phi_i' dx}{\int_0^L \rho \phi \phi_i dx}$$

$$\int_0^L (u_x E (u_x)'' + T_i u_x u_x') dx = \frac{EI \omega_i^2}{\rho}$$

Using orthogonal property $\phi = \sum_{i=1}^{\infty} c_i \phi_i$

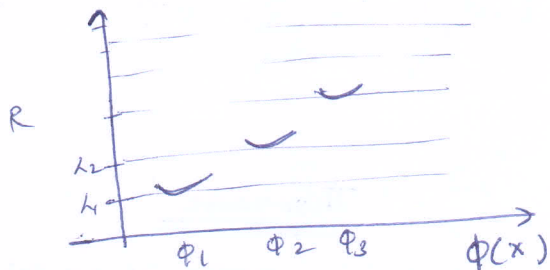
R = Rayleigh Quotient

$$\lambda_i = \frac{\int_0^L E \phi_i'^2 dx}{\int_0^L \rho \phi_i^2 dx}$$

$$\lambda = \min(R) = \frac{\int_0^L E \phi'^2 dx}{\int_0^L \rho \phi^2 dx}$$

A = Constant.

When ϕ become a particular mode shape then corresponding R or λ is local minimum (natural frequency)



$$\phi = \sum_{j=1}^{\infty} c_j \phi_j = c_1 \phi_1 + \sum_{j=1}^{\infty} \Delta c_j \phi_j$$

$$R = \frac{\int_0^L E (c_1 \phi_1')^2 dx + \int_0^L E (\Delta c_j \phi_j')^2 dx}{\int_0^L \rho (c_1 \phi_1)^2 dx + \int_0^L \rho (\Delta c_j \phi_j)^2 dx}$$

Recap

= free vibration of a bar
 - eigenvalue problem using calculus of variations

Rayleigh quotient

$$u(x, t) = \phi(x) q(t)$$

$$\phi(x) = \sum_{i=1}^{\infty} c_i \phi_i(x)$$

When A is not constant

$$R = \frac{\int_0^L EA \phi'^2 dx}{\int_0^L \rho A \phi^2 dx}$$

$$\left[\left(\frac{\partial u}{\partial x} \right)^2 + \left(\frac{\partial u}{\partial y} \right)^2 \right]$$

Minimum characterization of eigenvalue problem

[Sturm - Liouville problems]

$$\text{Min}_{\phi(x)} \int_0^L E \phi'^2 dx = \lambda_k$$

subject to

$$\int_0^L p \phi^2 dx - 1 = 0$$

$$\int_0^L p \phi_j \phi dx = 0$$

$$E \phi_i'' + \lambda_i p \phi_i = 0 \quad i = 1, 2, \dots, \infty$$

Theorem

If given constraint are satisfied then min value of obj fn will give eigenvalue λ_k .

(Normalize $\phi = 1$)

$j = 1, 2, 3, \dots, (k-1)$

(ϕ should be orthogonal to ϕ_j)

ϕ_j can be state variables

Governing eqⁿ

for ϕ_j (Eigen Value)

$$\phi = \sum_{i=1}^{\infty} c_i \phi_i$$

$$\int_0^L p \phi^2 dx = \sum_{i=1}^{\infty} c_i^2 = 1$$

$$\int_0^L p \phi \phi_j dx = 0 \Rightarrow \underline{c_j = 0}$$

$$E \phi_i'' + \lambda_i p \phi_i = 0 \quad \text{--- (3)}$$

Minimize

$$\text{Quantity} \int_0^L E \phi'^2 dx = \int_0^L E \phi' \left(\sum_{i=k}^{\infty} c_i \phi_i' \right) dx$$

$i = 1, \dots, (k-1)$

$$\int_0^L \phi_i \phi = 0$$

$$\int_0^L E \phi'^2 dx = \sum_{i=k}^{\infty} c_i \int_0^L E \phi' \phi_i' dx$$

Substitute from (3)

Integration by parts

$$\sum_{i=k}^{\infty} c_i \left[E \phi \phi_i' \Big|_0^L - \int_0^L E \phi \phi_i'' dx \right]$$

$$= \sum_{i=k}^{\infty} c_i \int_0^L h_i \rho \phi \phi_i dx$$

$$\sum_{i=k}^{\infty} h_i c_i^2$$

Eigen value problem convert a continuous function into discrete function.

Now we want to prove then $\min \sum_{i=k}^{\infty} h_i c_i^2 = h_k$

$$= \sum_{i=k}^{\infty} h_i c_i^2 + h_k \left(\underbrace{\sum_{i=k}^{\infty} c_i^2}_{=1 \text{ (Kaptos 1)}} \right) - h_k \left(\underbrace{\sum_{i=k}^{\infty} c_i^2}_{=1} \right)$$

$$\int_0^L E \phi^2 dx = h_k + \sum_{i=k}^{\infty} \underbrace{(h_i - h_k)}_{\geq 0} c_i^2$$

h_i are ordered
 $h_1 < h_2 < \dots < h_k < h_{k+1}$

$$\min \int_0^L E \phi^2 dx = h_k$$

Q for free vibration of a beam

$$R = \frac{\int_0^L EI \phi''^2 dx}{\int_0^L \rho \phi^2 dx}$$

When we $\min R$, we get h_k

See how to find Rayleigh Quotient for diff cases

Max natural frequency of a bar

$$\text{Max}_{A(x)} \quad \lambda = \frac{\int_0^L EA \phi'^2 dx}{\int_0^L \rho A \phi^2 dx}$$

R

subject to

$$\Delta: \int_0^L A dx - V^* \leq 0$$

Data: E, L, ρ, V^*

$$L = \frac{\int_0^L EA \phi'^2 dx}{\int_0^L \rho A \phi^2 dx} + \Delta \left(\int_0^L A dx - V^* \right)$$

$$\delta_\phi L = 0$$

$$\Rightarrow \frac{\left(\int_0^L \rho A \phi^2 dx \right) \left(\int_0^L 2EA \phi' \delta \phi' dx \right) - \left(\int_0^L EA \phi'^2 dx \right) \left(\int_0^L 2\rho A \phi \delta \phi dx \right)}{\left(\int_0^L \rho A \phi^2 dx \right)^2} > 0$$

$$\left(\int_0^L EA \phi' \delta \phi' dx \right) - \lambda \left(\int_0^L \rho A \phi \delta \phi dx \right) = 0$$

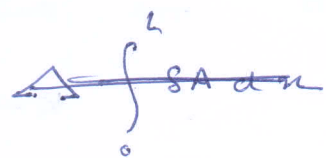
$$EA \phi' \delta \phi \Big|_0^L - \int_0^L (EA \phi'') \delta \phi dx - \lambda \int_0^L \rho A \phi \delta \phi dx = 0$$

$\delta\phi \rightarrow$ arbitrary

$$\boxed{(EA\phi')' + \lambda PA\phi = 0}$$

$\delta_A L = 0$

$$\left(\int_0^L PA\phi^2 dx \right) \left(\int_0^L EA\phi'^2 SA dx \right) - \left(\int_0^L EA\phi'^2 dx \right) \left(\int_0^L P\phi^2 SA dx \right) + \left(\int_0^L PA\phi^2 dx \right)^2 + \Delta \int_0^L SA dx = 0$$

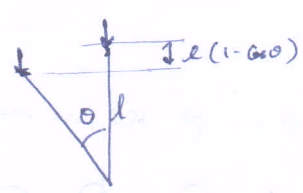
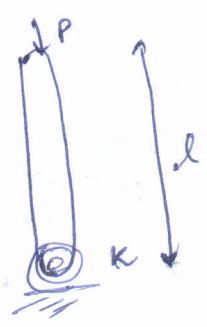
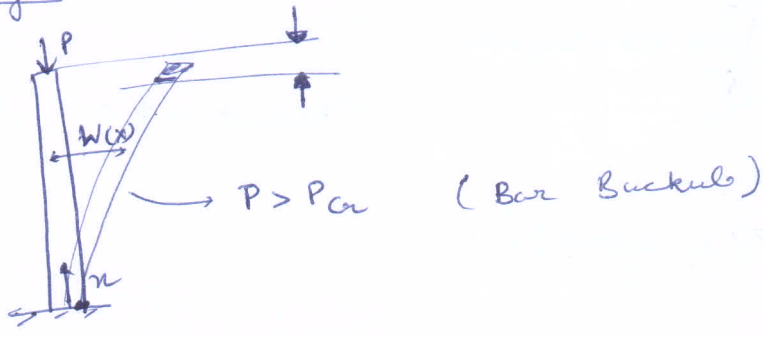


$$\left(\int_0^L EA\phi'^2 SA dx \right) - \lambda \left(\int_0^L P\phi^2 SA dx \right) + \Delta \left(\int_0^L PA\phi^2 dx \right) = 0$$

$$\boxed{EA\phi'^2 - \lambda P\phi^2 + \Delta \left(\int_0^L PA\phi^2 dx \right) = 0}$$

Design Equation

Buckling

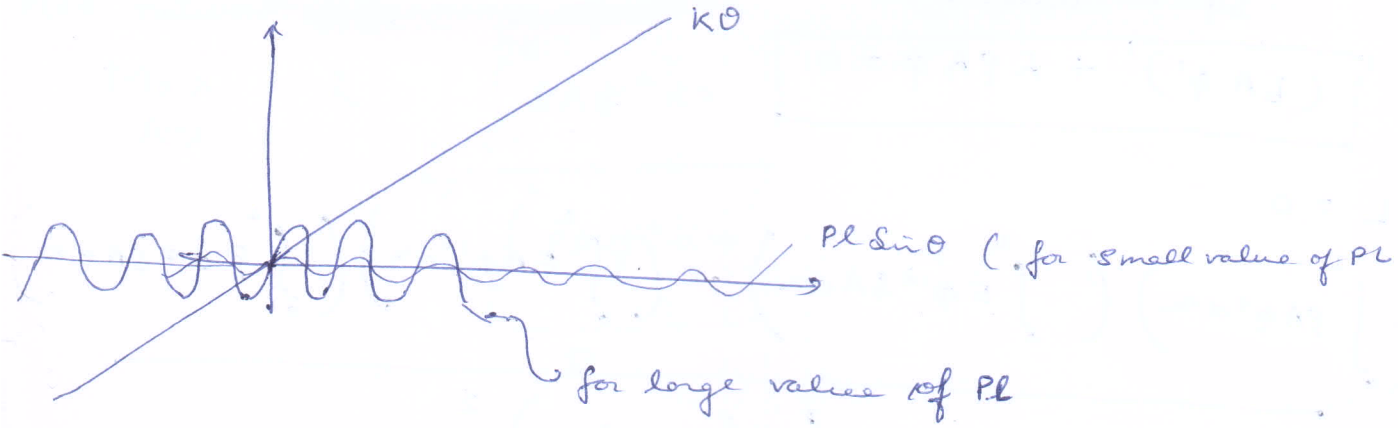


$$PE = \frac{1}{2} k\theta^2 + \left\{ -Pl(1 - \cos\theta) \right\}$$

for equilibrium

$$\frac{\partial PE}{\partial \theta} = 0 \Rightarrow k\theta - pl \sin\theta = 0 \quad \text{--- (1)}$$

$$\frac{\partial^2 PE}{\partial \theta^2} = k - pl \cos\theta \quad \text{--- (2)}$$

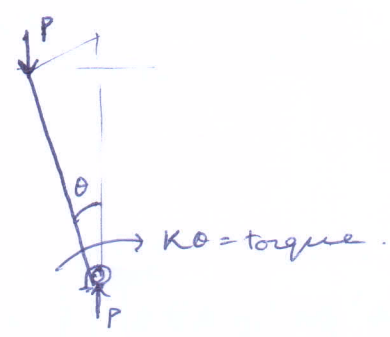
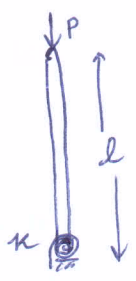


$$\frac{\partial^2 PE}{\partial \theta^2} > 0 \Rightarrow \text{min PE so all solution}$$

$$k - PL \cos \theta > 0$$

for $P > P_{cr}$ equilibrium straight position is no longer stable

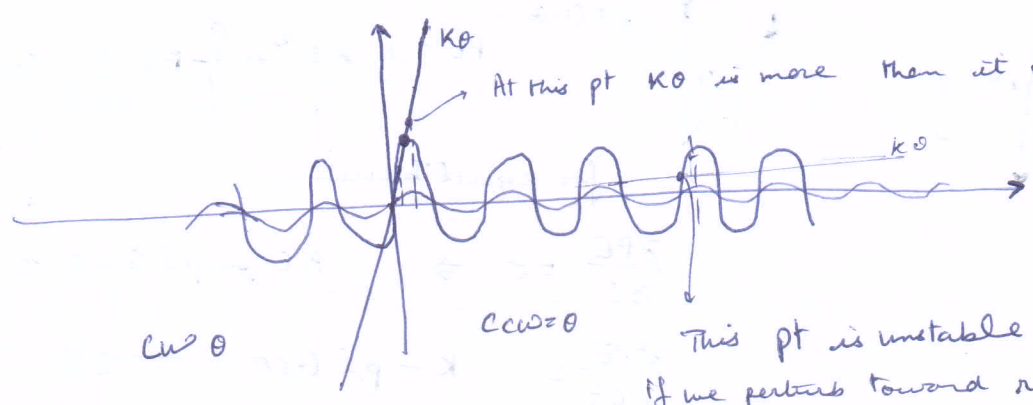
Buckling Stability



$$PL \sin \theta = k \theta \quad \text{moment balance}$$

$$PE = \frac{1}{2} k \theta^2 + \{ -P(l - l \cos \theta) \}$$

$$\frac{\partial PE}{\partial \theta} = k \theta - PL \sin \theta = 0 \quad (\text{Moment balance})$$



At this pt $k\theta$ is more than it means it is perturbed in CCW dir resulting our perturbation

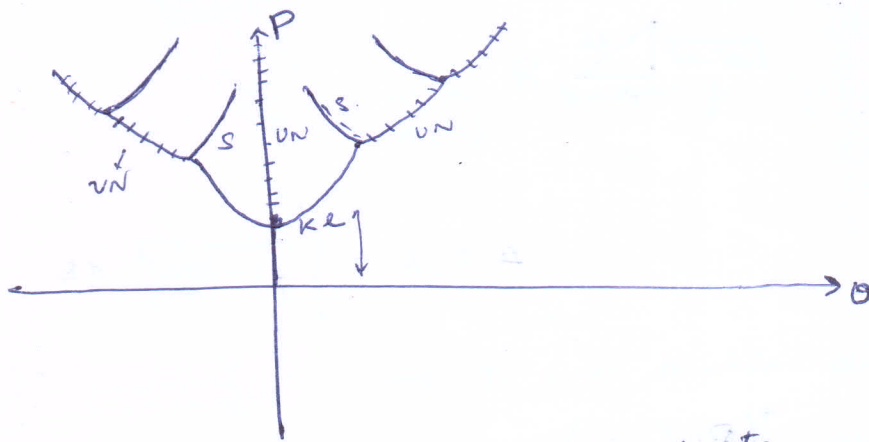
This pt is unstable if we perturb toward it $PL \sin \theta > k\theta$

$$\frac{\partial^2 PE}{\partial \theta^2} > 0$$

$$k - PL \cos \theta > 0$$

$$@ \theta = 0 \quad k - PL > 0$$

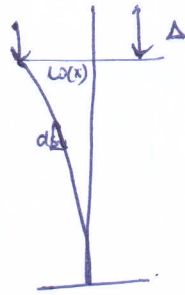
$$\frac{k}{L} > P \Rightarrow \text{min PE} \Rightarrow \text{stable sol}^n$$



Try to plot this in matlab.

Bifurcation in dynamic system

Buckling



$$\sqrt{ds^2 - dw^2} \quad \begin{array}{c} ds \\ dw \end{array}$$

$$PE = SE + WP$$

$$\Delta = L - \int_0^L \sqrt{1 - \left(\frac{dw}{ds}\right)^2} ds$$

$$PE = \int_0^L \frac{1}{2} EI w''^2 dx - P \Delta \quad \Delta = L - \int_0^L \left[1 - \frac{1}{2} \left(\frac{dw}{ds}\right)^2 \right] ds$$

$$\Delta = \int_0^L \frac{1}{2} \left(\frac{dw}{ds}\right)^2 ds$$

$$PE = \int_0^L \frac{1}{2} EI w''^2 dx - P \int_0^L \frac{1}{2} w'^2 dx$$

$$PE = \int_0^L \left[\frac{EI w''^2}{2} - \frac{P w'^2}{2} \right] dx$$

$$\delta_\omega PE = -(-Pw')' + (EIw'')'' = 0$$

$$(EIw'')'' + Pw'' = 0$$

This is twice differentiated moment eq.

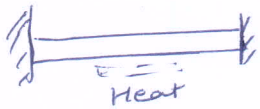
$$(EIw'')'' \rightarrow w''''$$

Buckling is a eigen value problem

Has infinite no of solution.

$$R_{\text{buckling}} = P = \frac{\int_0^L EI \phi''^2 dx}{\int_0^L \phi'^2 dx}$$

$$P_{cr} = \min_{\phi(x)} R_{\text{buckling}}$$



→ It will buckle due to residual comp stresses.

Most stable column

Max
 $A(x)$

Min
 $\phi(x)$

$$\frac{\int_0^L E \alpha A \phi''^2 dx}{\int_0^L \phi'^2 dx}$$

P_{cr}

Design of stable column

subject to

$$\int_0^L A dx - V^* \leq 0$$

Bower (Flat plate)

Recap

Buckling of columns in the framework of calculus of variations

Most stable column for given volume ($\max_{A(x)} \min_{\phi(x)} R_{\text{buckling}}$)

Code

Bar

$$Eu^2 = \Delta$$

$$A^{(k+1)} = A^{(k)} + (Eu^2 - \Delta)$$

or

$$A^{(k+1)} = A^{(k)} \left(\frac{Eu^2}{\Delta^{(k)}} \right)^{\eta}$$

$$\sum_{i=1}^n A_i^{(k+1)} l_i = V^*$$

Δ

$$\int A(x) dx - V^* \leq 0$$

$$\sum_{i=1}^n A_i^{(k)} l_i \left(\frac{Eu^2}{\Delta} \right)^{\eta} = V^*$$

$$\sum A_i l_i = V^*$$

$$\Delta^{\eta} V^* = \sum_{i=1}^n A_i^{(k)} l_i (Eu^2)^{\eta}$$

$$A^{(k+1)} = A^{(k)} \left(\frac{Eu^2}{\Delta} \right)^{\eta}$$

$$\sum A_i^{(k)} c_i^{\eta} l_i = V^*$$

$$\Delta = \left[\frac{\sum_{i=1}^n A_i^{(k)} l_i (Eu^2)^{\eta}}{V^*} \right]^{\frac{1}{\eta}}$$

$$A^{(k+1)} = A^{(k)} \left(\frac{Eu^2}{\Delta} \right)^{\eta}$$

$$\sum A_i l_i = A_{\max} l_i (\text{upcount}) + A_{\min} l_i (\text{downcount})$$

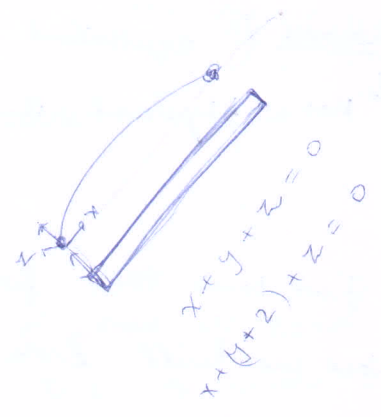
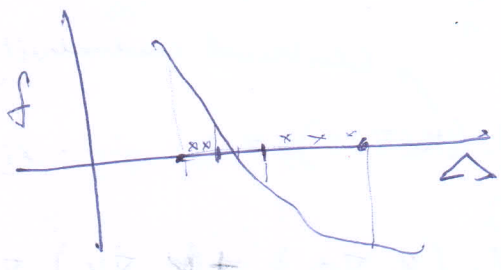
+ $\sum_{i \in R} A_i l_i$ \rightarrow These ones do not req μ_u & μ_d . The

These which have not reached upper or lower bound.

low for the remaining we are calculating lambda.
 for non linear variables, use bisection and other methods (iteratively)



In case of non linear fn.



zero for A_{max} & A_{min} , μ_u & $\mu_l \neq 0$
 but