KINEMATIC ANALYSIS OF THREE-LINK SPATIAL MECHANISMS CONTAINING SPHERE-PLANE AND SPHERE-GROOVE PAIRS

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Abstract—Kinematic pairs in a spatial mechanism are viewed either as allowing relative screw motion between links or as constraining the motion of the two chains of the mechanism connected to the two elements of the pair. Using pair geometry constraints of the sphere-plane and sphere-groove kinematic pairs, the displacement, velocity and acceleration equations are derived for, $R$-$Sp$-$R$, $R$-$Sp$-$P$, $P$-$Sp$-$P$, $P$-$Sp$-$R$ and $R$-$Sp$-$C$ three-link mechanisms. For known values of the input variable, other variables are computed in closed form. The analysis procedures are illustrated using numerical examples.

1. INTRODUCTION

The mechanisms containing higher pairs such as cams, sphere-plane, sphere-groove, or cylinder-plane provide the designer with the capabilities of designing machines and mechanisms to satisfy more complex and exact functional requirements than feasible with only lower pair mechanisms. These mechanisms in general are compact and contain fewer links than those with lower pairs.

In recent years, there has been considerable development in the tools for kinematic analysis of spatial mechanisms containing lower pairs.

Kinematic analysis of space mechanisms was initiated by the significant contribution of Dimentberg[1]. Dimentberg[2, 3] demonstrated the use of dual numbers and screw calculus to obtain closed-form displacement relationships of an RCCC* and other four-, five-, six- and seven-link spatial mechanisms containing revolute, cylinder, prismatic and helical pairs. Denavit[4] derived closed-form displacement relationships for a spatial RCCC mechanism using dual Euler angles. Yang[5] also derived such relationships for RCCC mechanisms using dual quaternions.

Vectors were first used by Chace[6] to derive closed-form displacement relations of RCCC mechanisms. Wallace and Freudenstein[7] also used vectors to obtain closed-form displacement relations of RRSRR and RRFRRR mechanisms.

Yang[8] proposed a general formulation using dual numbers to conduct displacement analysis of RCCRC spatial five-link mechanisms. Soni and Pamidi[9] extended this application of $(3 \times 3)$ matrices with dual elements to obtain closed-form displacement relations of RCCRR mechanisms.


Bagci[19] used a $(3 \times 3)$ screw matrix for displacement analysis of a mechanism containing two revo-
lute pairs, one cylinder pair and one spheric pair. Dobrovolski [20] used the method of spherical images to analyze space mechanisms containing revolute and cylinder pairs. Duffy [21, 22], Duffy and Habib–Olahi [23] used the method of spherical triangles to derive displacement relations for five and six link mechanisms containing revolute and cylinder pairs. Keller [25] and Gupta [26] also analyzed space mechanisms containing revolute, prismatic, cylinder, helical and spheric pairs. Recently Kohli and Soni [26] and Singh and Kohli [27] used the method of pair constraint geometry and successive screw displacements to conduct analyses of single and multi-loop mechanisms.

In the present paper, screw displacements expressed in vector form and the pair geometry constraints, also expressed in vector form, are used to derive the displacement, velocity and acceleration equations for \( R-Sp-R \), \( R-Sp-P \), \( P-Sp-R \), \( P-Sp-P \) and \( R-Sg-C \) three link mechanisms.

Since Revolute (R) and Prismatic (P) pairs are special cases of the cylinder pair (in prismatic pairs, the rotation is zero; for revolute pairs sliding is zero), we derive the analysis equation for \( C-Sp-C \) and \( C-Sg-C \) mechanisms, and then force rotations or translations at one or more pairs to zero, to obtain the equations for the above described three-link one degree of freedom mechanisms.

Briefly, the procedure for obtaining the analysis equations is as follows.

**Step 1.** Consider the \( C-Sp-C \) mechanism and the \( C-Sg-C \) mechanism.

**Step 2.** Separate the two moving links (Bodies 1 & 2) at the sphere–plane pair for the \( C-Sp-C \) case and at the sphere–groove pair for the \( C-Sg-C \).

**Step 3.** Use the screw displacements in vector form to describe the new \((j)\)th position of the sphere-plane \((Sp)\) or sphere-groove \((Sg)\) pairs from two sides of the pair.

**Step 4.** Use the pair geometry constraints on the position of the pair obtained from two sides.

**Step 5.** Force the cylindrical \((C)\) joints as revolute \((R)\) or prismatic \((P)\) joints by setting the sliding or the rotation equal to zero at cylindrical pairs.

### 2. THE THREE-LINK MECHANISM AND ASSOCIATED VECTORS

Figure 1 shows the initial position of two rigid bodies grounded via cylindrical pairs and connected together by a sphere–plane pair. Also shown are the following vectors and scalar quantities:

- \( u_a \) unit vector defining the direction of the axis of cylindric pair A.
- \( u_b \) unit vector defining the direction of the axis of cylindric pair B.
- \( P \) vector locating the axis of cylindric pair at A in the fixed coordinate system.
- \( Q \) vector locating the axis of cylindric pair at B in the fixed coordinate system.
- \( A \) unit vector perpendicular to the plane of the Sp pair embedded in body 1.
- \( A' \) vector embedded in body 2, congruent with A in the starting position, as shown in Fig. 1.
- \( R \) vector locating point R, the sphere center in the fixed coordinate system.
- \( \theta \) rotation of link 1 about axis \( u_a \).
- \( \theta_b \) rotation of link 2 about axis \( u_a \).
- \( S_a \) translation of link 1 along axis \( u_a \).
- \( S_b \) translation of link 2 along axis \( u_b \).

Figure 2 shows the \( C-Sg-C \) mechanism with all associated vectors and scalars. Description of all parameters are the same as for the \( C-Sp-C \) mechanism except for the direction of the vector A, which is now along the direction of the groove and also the addition of \( S \), which is the translation of the sphere along the direction of A.

### 3. PAIR GEOMETRY CONSTRAINT EQUATIONS

Figures 3 and 4 show a sphere–plane \((Sp)\) pair and a sphere–groove \((Sg)\) pair with the vector R locating R, the sphere center. The vector A, in the Sp pair is defined as a vector perpendicular to the plane in which the sphere moves. In the Sg pair, the vector A defines the direction of the groove.

![Fig. 1. C–Sp–C mechanism.](image1.png)

![Fig. 2. C–Sg–C mechanism.](image2.png)
We can now define the vectors $\mathbf{R}_j, \mathbf{A}_j, \mathbf{R}'_j$ and $\mathbf{A}'_j$. These new vectors will define the displaced position and direction of initially coincident point $R$ and vector $A$ in bodies 1 and 2 respectively after some relative motion between bodies 1 and 2. The prime notation here is used for new position expressed from the motion of body 2, whereas the unprimed notations are used for new positions expressed from the motion of body 1.

The pair geometry constraint equation for the $Sp$ pair is

$$\frac{d^n}{dt^n}([\mathbf{R}_j - \mathbf{R}'_j] \cdot \mathbf{A}_j) = 0, \quad n = 0, 1, 2, \ldots$$

which expresses that any relative motion between the sphere and the plane must be perpendicular to the vector $\mathbf{A}_j$ (Fig. 1).

The pair geometry constraint equation for the $Sg$ pair is

$$\frac{d^n}{dt^n} \mathbf{R}_j - \frac{d^n}{dt^n} \mathbf{R}'_j = \frac{d^n}{dt^n} (S_{oj}) \quad n = 0, 1, 2, \ldots$$

where $S_{oj}$ is the translation of the sphere along the groove in the direction of $\mathbf{A}'_j$. The constraint equation for the $Sg$ pair expresses that any relative motion between the sphere and the groove must be along the groove which is in the direction of $\mathbf{A}'_j$ (Fig. 2).

4. WORKING EQUATIONS

Referring to Fig. 1, let $\mathbf{A}$ be a vector in body 1 and $\mathbf{A}'$ a momentarily congruent vector in body 2 in the first position, perpendicular to the plane of the $Sp$

Fig. 3. Sphere-plane ($Sp$) pair.

pair. After some displacement of the mechanism, these vectors, in general, will separate due to the relative motion of the joint elements. Noting that both bodies 1 and 2 are connected to ground by $C$ pairs, we use the equations developed by Kohli and Soni[26] for expressing the direction of a vector embedded in the rigid body and also the displaced position of a point of the body after a rotation $\theta$ about the cylinder axis and a translation $S$ along the same axis. Using the prime notation for positions of the vector $\mathbf{A}'$ obtained from the motion of body 2 and the unprimed notation for positions of vector $\mathbf{A}$ (assumed frozen in body 1 in the first position and then moving with body 1) from the motion of body 1, the displaced directions of the vector $\mathbf{A}$ in bodies 1 and 2 are

$$\mathbf{A}_j = \cos \theta \mathbf{A} + \sin \theta (\mathbf{A} \times \mathbf{u}_a) + \mathbf{u}_a S_{aj} + \mathbf{P} \quad (3)$$

$$\mathbf{A}'_j = \cos \theta \mathbf{A}' + \sin \theta (\mathbf{A}' \times \mathbf{u}_a) + \mathbf{u}_a S_{aj}$$

Also, the displaced position of the point $R$ in rigid bodies 1 and 2 are given by:

$$\mathbf{R}_j = \cos \theta \mathbf{R} + \sin \theta (\mathbf{R} \times \mathbf{u}_a) + \mathbf{u}_a S_{aj} + \mathbf{P} \quad (5)$$

$$\mathbf{R}'_j = \cos \theta \mathbf{R} - \sin \theta (\mathbf{R} \times \mathbf{u}_a) + \mathbf{u}_a S_{aj} + \mathbf{Q} \quad (6)$$

Using the identity $[\mathbf{A} - (\mathbf{A} \times \mathbf{u}_a)] = (\mathbf{u}_a \times \mathbf{A}) + \mathbf{u}_a$, introducing the vectors

$$\mathbf{K} = \mathbf{R} - \mathbf{P} \quad (7)$$

$$\mathbf{L} = \mathbf{R} - \mathbf{Q} \quad (8)$$

and the following notation for any two vectors $\mathbf{u}_c$ and $\mathbf{D}$,

$$\mathbf{U}_{cd} = (\mathbf{u}_c \times \mathbf{D}) \times \mathbf{u}_c \quad (7a)$$

we can substitute eqns (7) and (7a) into eqns (5) and (6) to get

$$\mathbf{R}_j = \mathbf{R} + \mathbf{u}_a S_{aj} + (\cos \theta - 1) \mathbf{U}_{ak} + \sin \theta (\mathbf{u}_a \times \mathbf{K}) \quad (5a)$$

and

$$\mathbf{R}'_j = \mathbf{R} + \mathbf{u}_a S_{aj} + (\cos \theta - 1) \mathbf{U}_{al} + \sin \theta (\mathbf{u}_a \times \mathbf{L}) \quad (6a)$$

We now take the time-derivatives of equations for $\mathbf{R}_j$ and $\mathbf{R}'_j$ and using the notation of dots above the
variables to indicate time derivatives, we obtain the following equations

\[ \mathbf{R}_i = u_i \mathbf{S}_A + \left[ \cos \theta_{ij} (u_i \times \mathbf{K}) - \sin \theta_{ij} (u_i \times \mathbf{A}) \right] \theta_{ij} \]  

\[ \mathbf{R}'_i = u_i \mathbf{S}_B + \left[ \cos \theta_{ij} (u_i \times \mathbf{L}) - \sin \theta_{ij} (u_i \times \mathbf{B}) \right] \theta_{ij} \]  

\[ \mathbf{R}_i = u_i \mathbf{S}_A + \left[ \cos \theta_{ij} (u_i \times \mathbf{K}) + \sin \theta_{ij} (u_i \times \mathbf{A}) \right] \mathbf{B}_{ij} \]

\[ \mathbf{R}'_i = u_i \mathbf{S}_B + \left[ \cos \theta_{ij} (u_i \times \mathbf{L}) + \sin \theta_{ij} (u_i \times \mathbf{B}) \right] \mathbf{B}_{ij} \]  

Substituting eqn (7a) into eqn (4), using eqns (5a) and (6a), and by making the following substitutions

\[ M_{ij} = \cos \theta_{ij} (u_i \times \mathbf{K}) - \sin \theta_{ij} (u_i \times \mathbf{A}) \]  

\[ M_{ij} = \cos \theta_{ij} (u_i \times \mathbf{L}) - \sin \theta_{ij} (u_i \times \mathbf{B}) \]

\[ N_{ij} = \cos \theta_{ij} (u_i \times \mathbf{K}) + \sin \theta_{ij} (u_i \times \mathbf{A}) \]  

\[ N_{ij} = \cos \theta_{ij} (u_i \times \mathbf{L}) + \sin \theta_{ij} (u_i \times \mathbf{B}) \]  

we can derive the following working equations

\[ \mathbf{A}_i = A + (\cos \theta_{ij} - 1)u_i \mathbf{A} + \sin \theta_{ij} (u_i \times \mathbf{A}) \]  

\[ \mathbf{A}'_i = \left[ \cos \theta_{ij} (u_i \times \mathbf{A}) - \sin \theta_{ij} (u_i \times \mathbf{B}) \right] \mathbf{B}_{ij} \]  

\[ \mathbf{A}_i = \left[ \cos \theta_{ij} (u_i \times \mathbf{A}) - \sin \theta_{ij} (u_i \times \mathbf{B}) \right] \mathbf{B}_{ij} \]

\[ \mathbf{A}'_i = \left[ \cos \theta_{ij} (u_i \times \mathbf{A}) + \sin \theta_{ij} (u_i \times \mathbf{B}) \right] \mathbf{B}_{ij} \]

for the \( S_p \) pair.

\[ (\mathbf{R}_i - \mathbf{R}'_i) \cdot \mathbf{A}_i = 0. \]  

(20)

for the \( S_g \) pair.

\[ (\mathbf{R}_i - \mathbf{R}'_i) = \mathbf{A}_i S_{ij}. \]  

(21)

Observe that eqns (20) and (21) are eqns (1) and (2) with \( n = 0 \).

The cylindrical pairs used in the derivation may be forced to work as prismatic (P) pairs by letting \( \theta = 0 \) or may be forced to work as revolute (R) pairs by letting \( S = 0 \).

5.1 The P–Sp–P case

For this mechanism, we use \( \theta = \theta_b = 0 \) and eqns (14) and (17) are simplified to

\[ \mathbf{R}_i - \mathbf{R}'_i = u_i S_{Aij} - u_b S_{Bij} \]

and

\[ \mathbf{A}_i = \mathbf{A}. \]

Substituting in eqn (20), we get

\[ (u_i S_{Aij} - u_b S_{Bij}) \cdot \mathbf{A} = 0 \]

(22)

which simplifies to the input/output equation

\[ S_{Bij} = \frac{u_i \cdot \mathbf{A}}{u_p \cdot \mathbf{A}} S_{Aij}. \]  

(23)

5.2 The R–Sp–P case

\( \theta \) is the input; \( S_a \) is the output and \( \theta_b \equiv S_a \equiv 0 \). Equations (14) and (17) with \( \theta_b = S_a \equiv 0 \) substituted in eqn (20) provide,

\[ [- u_b S_{Bij} + (\cos \theta_{ij} - 1)u_i \mathbf{A} + \sin \theta_{ij} (u_i \times \mathbf{K})] \cdot \mathbf{A} = 0. \]

After simplification we obtain

\[ S_{Bij} = \frac{[(\cos \theta_{ij} - 1)u_i \mathbf{A} + \sin \theta_{ij} (u_i \times \mathbf{K})] \cdot \mathbf{A}}{u_p \cdot \mathbf{A}}. \]  

(24)

5.3 The R–Sp–R case

We have for this case \( S_a \equiv S_p \equiv 0 \), and eqn (14) and (17) are simplified to obtain

\[ \mathbf{R}_i - \mathbf{R}'_i = (\cos \theta_{ij} - 1)u_i \mathbf{A} + \sin \theta_{ij} (u_i \times \mathbf{K}) \]  

and

\[ \mathbf{A}_i = A + (\cos \theta_{ij} - 1)u_i \mathbf{A} + \sin \theta_{ij} (u_i \times \mathbf{A}) \]

Substituting the above equations into eqn (20), and simplifying the resulting equation, we obtain

\[ - S \cdot \mathbf{A} + (\cos \theta_{ij} - 1)(- u_p \cdot \mathbf{A} - S \cdot u_i \mathbf{A}) \]

\[ + \sin \theta_{ij} (u_i \times \mathbf{B}) \cdot \mathbf{A} - S \cdot (u_i \times \mathbf{A}) = 0 \]  

(25)
where \( \mathbf{S}_j \) is the known vector
\[
\mathbf{S}_j = (\cos \theta_{ij} - 1)\mathbf{U}_{ik} + \sin \theta_{ij}(\mathbf{u}_i \times \mathbf{K}).
\]  
(26)

Equation (25) can be solved for \( \theta_{ij} \) by using the following identities
\[
\cos \theta_{ij} = \frac{1 - \tan^2 \frac{\theta_{ij}}{2}}{1 + \tan^2 \frac{\theta_{ij}}{2}}; \quad \sin \theta_{ij} = \frac{2 \tan \frac{\theta_{ij}}{2}}{1 + \tan^2 \frac{\theta_{ij}}{2}}
\]
(27)
and simplifying the resulting quadratic equation to yield
\[
\tan \frac{\theta_{ij}}{2} = -\frac{b \pm \sqrt{(b^2 - c(c - 2a))}}{c - 2a}
\]
(28)

where:
\[
a = -\mathbf{U}_{ik} \cdot \mathbf{A} - \mathbf{S}_j \cdot \mathbf{U}_{ik} \\
b = (\mathbf{u}_i \times \mathbf{L}) \cdot \mathbf{A} - \mathbf{S}_j \cdot (\mathbf{u}_i \times \mathbf{A}) \\
c = -\mathbf{S}_j \cdot \mathbf{A}.
\]

5.4 The \( P-\text{Sp}-\text{R} \) case
Here, \( \theta_{ij} = S_{ij} = 0 \) and we have
\[
\mathbf{R}_j - \mathbf{R}_j' = \mathbf{u}_i S_{ij} - (\cos \theta_{ij} - 1)\mathbf{U}_{ik} - \sin \theta_{ij}(\mathbf{u}_i \times \mathbf{L})
\]
and
\[
\mathbf{A}'_j = \mathbf{A} + (\cos \theta_{ij} - 1)\mathbf{U}_{ik} + \sin \theta_{ij}(\mathbf{u}_i \times \mathbf{L}).
\]

Substituting the equations above into eqn (20) and simplifying, we get
\[
(\cos \theta_{ij} - 1)[-(\mathbf{U}_{ik} \cdot \mathbf{A} - \mathbf{S}_j \cdot (\mathbf{u}_i \times \mathbf{U}_{ik})] \\
+ \sin \theta_{ij}(\mathbf{u}_i \times \mathbf{L}) \cdot \mathbf{A} - \mathbf{S}_j \cdot (\mathbf{u}_i \times \mathbf{A}) - \mathbf{S}_j \mathbf{A}'_j \cdot \mathbf{A} = 0.
\]
(29)

Substituting quadratic in eqn (29) and simplifying the resulting quadratic gives us
\[
\tan \frac{\theta_{ij}}{2} = -\frac{b \pm \sqrt{(b^2 - c(c - 2a))}}{c - 2a}
\]
(30)

where this time
\[
a = -\mathbf{U}_{ik} \cdot \mathbf{A} - \mathbf{S}_j \cdot (\mathbf{u}_i \times \mathbf{U}_{ik}) \\
b = (\mathbf{u}_i \times \mathbf{L}) \cdot \mathbf{A} - \mathbf{S}_j \cdot (\mathbf{u}_i \times \mathbf{A}) \\
c = -\mathbf{S}_j \cdot \mathbf{A}.
\]

5.5 The \( R-\text{Sp}-\text{C} \) case
Only \( S_{ij} \) in eqn (17) is identically zero, so we get
\[
\mathbf{R}_j - \mathbf{R}_j' = -\mathbf{u}_i S_{ij} + \mathbf{S}_j - (\cos \theta_{ij} - 1)\mathbf{U}_{ik} \\
- \sin \theta_{ij}(\mathbf{u}_i \times \mathbf{L})
\]
where \( \mathbf{S}_j \) is given by eqn (26). Also,
\[
\mathbf{A}'_j = \mathbf{A} + (\cos \theta_{ij} - 1)\mathbf{U}_{ik} + \sin \theta_{ij}(\mathbf{u}_i \times \mathbf{A}).
\]

Substituting in eqn (21), we have
\[
\mathbf{u}_i S_{ij} - \mathbf{S}_j + (\cos \theta_{ij} - 1)\mathbf{U}_{ik} \\
+ \sin \theta_{ij}(\mathbf{u}_i \times \mathbf{L}) + \mathbf{A}'_j = 0.
\]

Forming the dot product of eqn (31) with \( \mathbf{A}'_j \times \mathbf{u}_i \) and upon simplification, we get
\[
\cos \theta_{ij}[S_j \cdot (\mathbf{A} \times \mathbf{u}_i) + \mathbf{U}_{ik} \cdot (\mathbf{A} \times \mathbf{u}_i)] \\
+ \sin \theta_{ij}[S_j \cdot \mathbf{U}_{ik} + \mathbf{U}_{ik} \cdot \mathbf{U}_{ik} - (\mathbf{u}_i \times \mathbf{L}) \cdot \mathbf{U}_{ik} = 0.
\]
(32)

Again, \( \theta_{ij} \) can be obtained by substituting eqns (27) into eqn (32) to obtain a quadratic whose solutions are
\[
\tan \frac{\theta_{ij}}{2} = -\frac{b \pm \sqrt{(a^2 + b^2 - c^2)}}{c - a}
\]
(33)

where
\[
a = S_j \cdot (\mathbf{A} \times \mathbf{u}_i) + \mathbf{U}_{ik} \cdot (\mathbf{A} \times \mathbf{u}_i) \\
b = S_j \cdot \mathbf{U}_{ik} + \mathbf{U}_{ik} \cdot \mathbf{U}_{ik} \\
c = - (\mathbf{u}_i \times \mathbf{L}) \cdot \mathbf{U}_{ik}.
\]

Forming the dot product of eqn (31) with \( \mathbf{u}_i \times \mathbf{L} \) and simplifying, we get
\[
S_{ij} = \frac{[S_j \cdot (\mathbf{u}_i \times \mathbf{L}) - (\cos \theta_{ij} - 1)\mathbf{U}_{ik} \cdot (\mathbf{u}_i \times \mathbf{L})]}{\mathbf{A}'_j \cdot (\mathbf{u}_i \times \mathbf{L})} - \frac{\sin \theta_{ij}(\mathbf{u}_i \times \mathbf{L}) \cdot (\mathbf{u}_i \times \mathbf{L})}{\mathbf{A}'_j \cdot (\mathbf{u}_i \times \mathbf{L})}.
\]
(34)

Forming the dot product of eqn (31) with \( \mathbf{u}_i \) and simplifying, we get
\[
S_{ij} = [S_j - (\cos \theta_{ij} - 1)\mathbf{U}_{ik} - \sin \theta_{ij}(\mathbf{u}_i \times \mathbf{L})] \\
- S_{ij} \mathbf{A}'_j \cdot \mathbf{u}_i.
\]
(35)

6. VELOCITY AND ACCELERATION ANALYSIS

To obtain the velocity and acceleration relations, we can either (a) take the derivatives with respect to time of the displacement equations or (b) use the higher order constraint equations. For the \( P-\text{Sp}-\text{P} \) case, taking the derivative of the displacement equation is trivial. But for the other cases, this procedure is cumbersome. It is therefore more convenient to just use eqns (14)–(19) in the following constraint eqns (36)–(39), which are eqns (1) and (2) with \( n = 1 \) and \( n = 2 \).

For the \( \text{Sp} \) pair
\[
(\mathbf{R}_j - \mathbf{R}_j') \cdot \mathbf{A}'_j + (\mathbf{R}_j - \mathbf{R}_j') \cdot \mathbf{A}'_j = 0
\]
(36)
and
\[
(\ddot{\mathbf{R}}_j - \ddot{\mathbf{R}}_j') \cdot \mathbf{A}'_j + 2(\dot{\mathbf{R}}_j - \dot{\mathbf{R}}_j') \cdot \mathbf{A}'_j + (\mathbf{R}_j - \mathbf{R}_j') \cdot \ddot{\mathbf{A}}_j = 0.
\]
(37)
For the \( S_g \) pair

\[
\dot{\mathbf{R}}_i - \dot{\mathbf{R}}_j = A_i' \dot{S}_{iG} + \dot{\mathbf{A}}_j' S_{jG} \tag{38}
\]

and

\[
\ddot{\mathbf{R}}_i - \ddot{\mathbf{R}}_j = \ddot{\mathbf{A}}_j S_{jG} + 2 \dot{\mathbf{A}}_j' \dot{S}_{jG} + \dot{\mathbf{A}}_j' S_{jG} \tag{39}
\]

6.1 The P-Sp-P case

Here we can use the time derivatives of the displacement equation to get

\[
\ddot{S}_{iG} = \frac{u_x \cdot A}{u_y \cdot A} \ddot{S}_{ai} \tag{40}
\]

\[
\ddot{S}_{jG} = \frac{u_x \cdot A}{u_y \cdot A} \ddot{S}_{aj} \tag{41}
\]

6.2 The R-Sp-P case

Equations (18) and (19) become

\[
\dot{\mathbf{R}}_i - \dot{\mathbf{R}}_j = M_{ij} \dot{\theta}_{aj} - u_y \dot{S}_{ij} \tag{42}
\]

and

\[
\dot{\mathbf{R}}_i - \dot{\mathbf{R}}_j = - N_{ij} \dot{\theta}_{aj} + M_{ij} \ddot{\theta}_{aj} - u_y \dot{S}_{ij} \tag{43}
\]

6.3 The R-Sp-R case

\[
S_{ai} = S_{aj} = \ddot{S}_{ai} = \ddot{S}_{aj} = \ddot{S}_{bi} = \ddot{S}_{bj} \equiv 0. \tag{44}
\]

Equations (18) and (9) become

\[
\dot{\mathbf{R}}_i - \dot{\mathbf{R}}_j = M_{ij} \dot{\theta}_{aj} - M_{ij} \dot{\theta}_{aj} \tag{45}
\]

and

\[
\dot{\mathbf{R}}_i - \dot{\mathbf{R}}_j = - N_{ij} \dot{\theta}_{aj} + M_{ij} \ddot{\theta}_{aj} + N_{ij} \ddot{\theta}_{aj} \tag{46}
\]

Also,

\[
\dot{A}_j = \dot{V}_j, \quad \dot{\dot{A}}_j = \dot{V}_j \times \dot{u}_j \tag{47}
\]

Substituting in eqns (36) and (37), we get

\[
\ddot{\mathbf{S}}_{ej} - u_x \cdot \dot{A}_j = \frac{u_x \cdot \dot{A}_j}{u_y \cdot \dot{A}_j} \ddot{\mathbf{S}}_{ej} \tag{48}
\]

Substituting the expression for (\( \dot{\mathbf{R}}_j - \dot{\mathbf{R}}_j \)) just obtained into eqn (38) we get

\[
M_{ij} \dot{\theta}_{aj} - u_y \dot{S}_{ij} = A_j' S_{ijG} + \dot{V}_{ij} S_{ijG} \tag{49}
\]

Substituting in eqns (36) and (37), we get

\[
\dot{\theta}_{aj} = \frac{M_{ij} \cdot \dot{A}_j}{M_{ij} \cdot \dot{A}_j - (\mathbf{R}_j - \mathbf{R}_j) \cdot \mathbf{V}_{ij}} \ddot{\mathbf{S}}_{ij} \tag{50}
\]

\[
\dot{\theta}_{aj} = \frac{- N_{ij} \cdot \dot{A}_j}{D} + \frac{M_{ij} \cdot \dot{A}_j}{D} + \frac{2M_{ij} \cdot \mathbf{V}_{ij} \cdot \dot{\theta}_{aj}}{D} + \frac{N_{ij} \cdot \dot{A}_j - 2M_{ij} \cdot \ddot{\mathbf{V}}_{ij} - (\mathbf{R}_j - \mathbf{R}_j) \cdot \mathbf{W}_{ij} \cdot \dot{\theta}_{aj}}{D} \tag{51}
\]

where

\[
D = M_{ij} \cdot \dot{A}_j - (\mathbf{R}_j - \mathbf{R}_j) \cdot \mathbf{V}_{ij} \tag{52}
\]

6.4 The P-Sp-R case

Equations (18) and (19) are

\[
\dot{\mathbf{R}}_i - \dot{\mathbf{R}}_j = u_x \dot{S}_{ai} - M_{ij} \dot{\theta}_{aj} \tag{53}
\]

\[
\ddot{\mathbf{R}}_i - \ddot{\mathbf{R}}_j = u_x \ddot{S}_{ai} + N_{ij} \ddot{\theta}_{aj} - M_{ij} \dot{\theta}_{aj} \tag{54}
\]

Also,

\[
\dot{A}_j = \mathbf{V}_j \ddot{\theta}_{aj} \tag{55}
\]

Substituting in eqns (36) and (37) we get

\[
\dot{\theta}_{aj} = \frac{u_x \cdot \dot{A}_j}{M_{ij} \cdot \dot{A}_j - (\mathbf{R}_j - \mathbf{R}_j) \cdot \mathbf{V}_{ij}} \ddot{\mathbf{S}}_{ij} \tag{56}
\]

6.5 The R-Sp-C case

Only \( S_{ai} \), \( S_{aj} \) and \( S_{bi} \) are zero and eqns (18) and (19) become:

\[
\dot{\mathbf{R}}_i - \dot{\mathbf{R}}_j = M_{ij} \dot{\theta}_{aj} - u_y \dot{S}_{ij} - M_{ij} \dot{\theta}_{aj} \tag{57}
\]

and

\[
\ddot{\mathbf{R}}_i - \ddot{\mathbf{R}}_j = - N_{ij} \dot{\theta}_{aj} + M_{ij} \ddot{\theta}_{aj} - u_y \dot{S}_{ij} + N_{ij} \ddot{\theta}_{aj} - M_{ij} \dot{\theta}_{aj} \tag{58}
\]

Also,

\[
A_j' = V_j' \dot{\theta}_{aj} \tag{59}
\]

Substituting the expression for \( \dot{\mathbf{R}}_j - \dot{\mathbf{R}}_j \) just obtained into eqn (38) we get

\[
(M_{ij} \dot{\theta}_{aj} - u_y \dot{S}_{ij} - M_{ij} \dot{\theta}_{aj}) - A_j' \dot{S}_{ij} + \mathbf{V}_j' S_{ij} = 0 \tag{60}
\]

Forming the dot product of eqn (49) with \( \dot{A}_j' \times u_y \), we get

\[
(M_{ij} \dot{\theta}_{aj} - u_y \dot{S}_{ij} - M_{ij} \dot{\theta}_{aj}) \cdot \dot{A}_j' \times u_y = V_j' \cdot (A_j' \times u_y) \dot{\theta}_{aj} S_{ij} \tag{61}
\]

or

\[
\dot{\theta}_{aj} = \frac{M_{ij} \cdot \dot{A}_j' \times u_y}{(S_{ij} \cdot \mathbf{V}_j' + M_{ij} \cdot \dot{A}_j' \times u_y)} \tag{62}
\]

Now, forming the dot product of eqn (49) with \( A_j' \times (S_{ij} \cdot \mathbf{V}_j' + M_{ij}) \), we have

\[
(M_{ij} \dot{\theta}_{aj} - u_y \dot{S}_{ij}) \cdot A_j' \times (S_{ij} \cdot \mathbf{V}_j' + M_{ij}) = 0 \tag{63}
\]
or

\[ \hat{S}_{ij} = \frac{M_{ij} \cdot \hat{A}_i \times (S_{ij} V_B + M_{ij})}{u_B \cdot \hat{A}_i \times (S_{ij} V_B + M_{ij})} \hat{A}_j \cdot \hat{A}_i \times \hat{A}_j \times (S_{ij} V_B + M_{ij}) = 0 \]  

(51)

Again forming the dot product of eqn (49) with

\[ u_B \times (S_{ij} V_B + M_{ij}), \]

we have

\[ (M_{ij} \hat{A}_j - \hat{A}_j \hat{A}_i \cdot u_B \times (S_{ij} V_B + M_{ij}) = 0 \]

or

\[ \hat{S}_{ij} = \frac{M_{ij} \cdot u_B \times (S_{ij} V_B + M_{ij})}{u_B \cdot \hat{A}_i \times (S_{ij} V_B + M_{ij})} \hat{A}_j \cdot \hat{A}_i \times \hat{A}_j \times (S_{ij} V_B + M_{ij}) = 0 \]  

(52)

Acceleration: Substituting the expression for \( \hat{R}_i \) obtained earlier for the \( R-Sp-R \) case into eqn (39), we will get

\[ -N_A \hat{\theta}_A + M_A \hat{\theta}_A - u_B \hat{S}_B + N_B \hat{\theta}_B - M_B \hat{\theta}_B = \hat{A}_j \hat{S}_{ij} + 2V_B \hat{\theta}_B \hat{S}_{ij} + A \hat{\theta}_j \hat{S}_{ij} \]

or

\[ u_B \hat{\theta}_B + \hat{A}_j \hat{S}_{ij} + (S_{ij} V_B + M_{ij}) \hat{\theta}_B = -N_A \hat{\theta}_A + M_A \hat{\theta}_A + (N_B + S_{ij} W_B) \hat{\theta}_B + 2V_B \hat{\theta}_B \hat{S}_{ij} + A \hat{\theta}_j \hat{S}_{ij} \]  

(53)

Letting \( X \) be equal to the r.h.s. of eqn (53) and by using the same technique of forming the dot product of eqn (53) with the proper cross-products, we will obtain the following

\[ \hat{\theta}_B = \frac{X \cdot \hat{A}_j \times u_B}{(S_{ij} V_B + M_{ij}) \cdot \hat{A}_j \times u_B} \]  

(54)

\[ \hat{S}_{ij} = \frac{X \cdot \hat{A}_j \times (S_{ij} V_B + M_{ij})}{u_B \cdot \hat{A}_i \times (S_{ij} V_B + M_{ij})} \]  

(55)

\[ \hat{\theta}_j = \frac{X \cdot u_B \times (S_{ij} V_B + M_{ij})}{\hat{A}_j \cdot u_B \times (S_{ij} V_B + M_{ij})} \]  

(56)

7. NUMERICAL EXAMPLES

1. Analysis of a \( R-Sp-R \) mechanism.
   The vectors describing the mechanism are

\[ \begin{align*}
   u_A &= 0i + 1j + 0k \\
   u_B &= (3i + 1j + 0k)/\sqrt{10} \\
   P &= 0i + 0j + 0k \\
   Q &= 0i + 4j + 0.75k \\
   R &= 1i + 1.5j + 2k \\
   A &= 0i + 0j + 1k
\end{align*} \]

The plot of the output displacement (\( \theta_B \)), velocity (\( \hat{\theta}_B \)) and acceleration (\( \hat{\theta}_B \)) are given in Fig. 5.

2. Displacement, velocity and acceleration analysis of a \( R-Sp-C \) mechanism.

Fig. 5. Plot of \( \theta_B, \hat{\theta}_B \) and \( \hat{\theta}_B \) for the \( R-Sp-R \) mechanism.
Table 1. Displacements, velocities and accelerations

<table>
<thead>
<tr>
<th>θ_A</th>
<th>θ_B</th>
<th>θ_G</th>
<th>4_A</th>
<th>4_B</th>
<th>6_A</th>
<th>6_B</th>
<th>6_G</th>
<th>8_A</th>
<th>8_B</th>
<th>8_G</th>
<th>10_A</th>
<th>10_B</th>
<th>10_G</th>
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<td>1.76</td>
<td>.36</td>
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</tr>
<tr>
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<td>12.99</td>
<td>.12</td>
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<td>-2.06</td>
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<td>5.38</td>
<td>2.04</td>
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<tr>
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<td>32.3</td>
<td>-2.06</td>
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<td>5.38</td>
<td>2.04</td>
<td>-.84</td>
<td>-4.07</td>
<td></td>
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<tr>
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<td>-.55</td>
<td>1.11</td>
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<tr>
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<td>.08</td>
<td>1.32</td>
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<td>-1.20</td>
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<td>.52</td>
<td></td>
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</tr>
</tbody>
</table>

The mechanism parameters are

\[ u_1 = (2j + 1k)/\sqrt{6} \]
\[ u_2 = (2j + 1k)/\sqrt{2} \]
\[ P = 0i + 0j + 0k \]
\[ Q = 0i + 0j + 1k \]
\[ R = 3j + 3j + 3k \]
\[ A = (2j + 1j + 2k)/\sqrt{6}. \]

The motion parameters are: \( \dot{\theta}_j \) is one unit of angular velocity and \( \ddot{\theta}_k \) is zero, both constant for \( j = 0, 1, 2, \ldots \).

The results of the analysis for the \( R-S_\text{g}-C \) mechanism are shown in Table 1. The direction of the rotations and linear motions are established using the right hand rule. Rotations are positive counterclockwise looking at the head of the unit vectors \( u_1 \) and \( u_2 \). Linear motions are positive when they are in the direction of the vectors they are associated with.

It is to be mentioned here also that although the quadratic equations gave two sets of solutions, only one set will define the motion of the mechanism. The other set of solutions are for those positions in which the mechanism has to be disassembled into the other possible configuration.

8. CONCLUSIONS

Displacements, velocities and accelerations have been derived for several three-link spatial mechanisms containing sphere–plane and sphere–groove pairs. The groove of the sphere–groove pair was assumed to be a cylindrical groove, resulting in straight line axis of the groove. However, a more generalized groove may be one whose axis is a spatial curve. The authors are working on developing analysis procedures for mechanisms containing such a generalized sphere–groove pair. The expected results of their work will be the subject of a forthcoming paper. Similarly, the authors also have the generalization of the sphere–plane pair in progress, in which the parallel planes of the pair are generalized to form equidistant curved surfaces.

Acknowledgements—The authors wish to acknowledge the support of this research under Grant No. DAAG29-81K-0125 at the University of Florida sponsored by the U.S. Army Research Office. The second author also wishes to acknowledge the support of the College of Engineering at the University of Wisconsin, Milwaukee.

REFERENCES


APPENDIX

1. Sphere–plane constraint equation

The complete displacement constraint equations of the Sp-pair are

\[ \mathbf{R}_j - \mathbf{R}_j' = S_p \mathbf{w}_p \]  

(a)

and

\[ \mathbf{u}_p \cdot \mathbf{A}_j' = 0 \]  

(b)

where \( \mathbf{w}_p \) is a unit vector in the plane of the Sp pair, perpendicular to \( \mathbf{A}_j \), and is in the direction of the relative motion of point \( R \) of body 1 with respect to the initially coincident point \( R' \) of body 2.

Derivatives of equations (a) and (b) with respect to time are taken to give the following velocity and acceleration constraint equations

Velocity:

\[ \dot{\mathbf{R}}_j - \dot{\mathbf{R}}_j' = \dot{S}_p \mathbf{w}_p + S_p \dot{\mathbf{w}}_p \]  

(c)

and

\[ \dot{\mathbf{w}}_p \cdot \mathbf{A}_j' + \dot{\mathbf{w}}_p \cdot \dot{\mathbf{A}}_j = 0 \]  

(d)

Acceleration:

\[ \ddot{\mathbf{R}}_j - \ddot{\mathbf{R}}_j' = \ddot{S}_p \mathbf{w}_p + 2 \dot{S}_p \dot{\mathbf{w}}_p + S_p \ddot{\mathbf{w}}_p \]  

(e)

and

\[ \ddot{\mathbf{w}}_p \cdot \mathbf{A}_j' + 2 \dot{\mathbf{w}}_p \cdot \dot{\mathbf{A}}_j + \dot{\mathbf{w}}_p \cdot \ddot{\mathbf{A}}_j = 0 \]  

(f)

The constraint eqns (a)–(f) are complete in the sense that all of the important variables in the motion of the joint elements are included. Also, the Coriolis component in the acceleration constraint eqn (f) is evident since \( \dot{\mathbf{A}}_j \) is a function of \( \ddot{\mathbf{R}}_j \).

2. Proof that \( \frac{d^n}{dt^n}[(\mathbf{R}_j - \mathbf{R}_i) \cdot \mathbf{A}_j] = 0, n = 0, 1, 2 \) satisfies the complete Sp pair constraint equation

Without loss of generality, we can let \( S_p = S_p \mathbf{w}_p \) and write the complete constraint equation as

\[ \frac{d^n}{dt^n} (\mathbf{R}_j - \mathbf{R}_i) = \frac{d^n}{dt^n} (S_p) \]  

(a)

and

\[ \frac{d^n}{dt^n} (S_p \cdot \mathbf{A}_j) = 0. \]  

(b)

Displacement: For \( n = 0 \), eqn (a) and (b) are

\[ (\mathbf{R}_j - \mathbf{R}_i) = S_p \]  

(c)

and

\[ S_p \cdot \mathbf{A}_j = 0. \]  

(d)

Forming the dot product of eqn (c) with \( \mathbf{A}_j \) gives us the displacement constraint equation for the Sp pair.

\[ (\mathbf{R}_j - \mathbf{R}_i) \cdot \mathbf{A}_j = 0. \]  

(e)

Velocity: With \( n = 1 \), eqns (a) and (b) will become

\[ \dot{\mathbf{R}}_j - \dot{\mathbf{R}}_j' = \dot{S}_p \]  

(f)

and

\[ \dot{S}_p \cdot \mathbf{A}_j = - S_p \cdot \dot{\mathbf{A}}_j. \]  

(g)

Taking the dot product of eqn (f) with \( \mathbf{A}_j \) gives us

\[ (\dot{\mathbf{R}}_j - \dot{\mathbf{R}}_j') \cdot \mathbf{A}_j = \dot{S}_p \cdot \mathbf{A}_j. \]  

(h)

Substituting eqn (g) into (h), we will have

\[ (\dot{\mathbf{R}}_j - \dot{\mathbf{R}}_j') \cdot \mathbf{A}_j = - S_p \cdot \dot{\mathbf{A}}_j. \]  

(i)

Equation (c) can now be substituted in eqn (f) to get

\[ (\dot{\mathbf{R}}_j - \dot{\mathbf{R}}_j') \cdot \mathbf{A}_j = -(\mathbf{R}_j - \mathbf{R}_i) \cdot \mathbf{A}_j \]  

or

\[ (\dot{\mathbf{R}}_j - \dot{\mathbf{R}}_j') \cdot \mathbf{A}_j + (\mathbf{R}_j - \mathbf{R}_i) \cdot \dot{\mathbf{A}}_j = 0 \]  

(j)

which is really

\[ \frac{d}{dt} [(\mathbf{R}_j - \mathbf{R}_i) \cdot \mathbf{A}_j] = 0. \]  

(k)

Acceleration: For \( n = 2 \), eqns (a) and (b) will be

\[ \ddot{\mathbf{R}}_j - \ddot{\mathbf{R}}_j' = \ddot{S}_p \]  

(l)
Forming the dot product of eqn (1) with $\mathbf{A}'$ and substituting

$$\mathbf{S}_p \cdot \mathbf{A}' = 2\mathbf{S}_p \cdot \mathbf{A}' + \mathbf{S}_p \cdot \mathbf{A}' = 0$$

from eqn (m), we will get

$$(\mathbf{R}_j - \mathbf{R}_j) \cdot \mathbf{A}' = -2\mathbf{S}_p \cdot \mathbf{A}' + \mathbf{S}_p \cdot \mathbf{A}'$$

Substituting eqns (c) and (f) into eqn (n) gives us

$$(\mathbf{R}_j - \mathbf{R}_j) \cdot \mathbf{A}' + 2(\mathbf{R}_j - \mathbf{R}_j) \cdot \mathbf{A}' + (\mathbf{R}_j - \mathbf{R}_j) \cdot \mathbf{A}' = 0$$

which is

$$\frac{d^2}{dt^2}[(\mathbf{R}_j - \mathbf{R}_j) \cdot \mathbf{A}] = 0.$$